

# LOW ENERGY SOLUTIONS FOR THE SEMICLASSICAL LIMIT OF SCHROEDINGER MAXWELL SYSTEMS

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**ABSTRACT.** We show that the number of solutions of Schrödinger Maxwell system on a smooth bounded domain  $\Omega \subset \mathbb{R}^3$ . depends on the topological properties of the domain. In particular we consider the Lusternik-Schnirelmann category and the Poincaré polynomial of the domain.

## Dedicated to our friend Bernhard

### 1. INTRODUCTION

Given real numbers  $q > 0$ ,  $\omega > 0$  we consider the following Schrödinger Maxwell system on a smooth bounded domain  $\Omega \subset \mathbb{R}^3$ .

$$(1) \quad \begin{cases} -\varepsilon^2 \Delta u + u + \omega uv = |u|^{p-2}u & \text{in } \Omega \\ -\Delta v = qu^2 & \text{in } \Omega \\ u, v = 0 & \text{on } \partial\Omega \end{cases}$$

This paper deals with the semiclassical limit of the system (1), i.e. it is concerned with the problem of finding solutions of (1) when the parameter  $\varepsilon$  is sufficiently small. This problem has some relevance for the understanding of a wide class of quantum phenomena. We are interested in the relation between the number of solutions of (1) and the topology of the bounded set  $\Omega$ . In particular we consider the Lusternik Schnirelmann category  $\text{cat } \Omega$  of  $\Omega$  in itself and its Poincaré polynomial  $P_t(\Omega)$ .

Our main results are the following.

**Theorem 1.** *Let  $4 < p < 6$ . For  $\varepsilon$  small enough there exist at least  $\text{cat}(\Omega)$  positive solutions of (1).*

**Theorem 2.** *Let  $4 < p < 6$ . Assume that for  $\varepsilon$  small enough all the solutions of problem (1) are non-degenerate. Then there are at least  $2P_1(\Omega) - 1$  positive solutions.*

Schrödinger Maxwell systems recently received considerable attention from the mathematical community. In the pioneering paper [9] Benci and Fortunato studied system (1) when  $\varepsilon = 1$  and without nonlinearity. Regarding the system in a semiclassical regime Ruiz [18] and D'Aprile-Wei [11] showed the existence of a family of radially symmetric solutions respectively for  $\Omega = \mathbb{R}^3$  or a ball. D'Aprile-Wei [12] also proved the existence of clustered solutions in the case of a bounded domain  $\Omega$  in  $\mathbb{R}^3$ .

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Recently, Siciliano [19] relates the number of solution with the topology of the set  $\Omega$  when  $\varepsilon = 1$ , and the nonlinearity is a pure power with exponent  $p$  close to the critical exponent 6. Moreover, in the case  $\varepsilon = 1$ , many authors proved results of existence and non existence of solution of (1) in presence of a pure power nonlinearity  $|u|^{p-2}u$ ,  $2 < p < 6$  or more general nonlinearities [1, 2, 3, 4, 10, 14, 15, 17, 20].

In a forthcoming paper [13], we aim to use our approach to give an estimate on the number of low energy solutions for Klein Gordon Maxwell systems on a Riemannian manifold in terms of the topology of the manifold and some information on the profile of the low energy solutions.

In the following we always assume  $4 < p < 6$ .

## 2. NOTATIONS AND DEFINITIONS

In the following we use the following notations.

- $B(x, r)$  is the ball in  $\mathbb{R}^3$  centered in  $x$  with radius  $r$ .
- The function  $U(x)$  is the unique positive spherically symmetric function in  $\mathbb{R}^3$  such that

$$-\Delta U + U = U^{p-1} \text{ in } \mathbb{R}^3$$

we remark that  $U$  and its first derivative decay exponentially at infinity.

- Given  $\varepsilon > 0$  we define  $U_\varepsilon(x) = U\left(\frac{x}{\varepsilon}\right)$ .
- We denote by  $\text{supp } \varphi$  the support of the function  $\varphi$ .
- We define

$$m_\infty = \inf_{\int_{\mathbb{R}^3} |\nabla v|^2 + v^2 dx = |v|_{L^p(\mathbb{R}^3)}^p} \frac{1}{2} \int_{\mathbb{R}^3} |\nabla v|^2 + v^2 dx - \frac{1}{p} |v|_{L^p(\mathbb{R}^3)}^p$$

- We also use the following notation for the different norms for  $u \in H_g^1(M)$ :

$$\begin{aligned} \|u\|_\varepsilon^2 &= \frac{1}{\varepsilon^3} \int_M \varepsilon^2 |\nabla u|^2 + u^2 dx & |u|_{\varepsilon,p}^p &= \frac{1}{\varepsilon^3} \int_\Omega |u|^p dx \\ \|u\|_{H_0^1}^2 &= \int_\Omega |\nabla u|^2 dx & |u|_p^p &= \int_\Omega |u|^p dx \end{aligned}$$

and we denote by  $H_\varepsilon$  the Hilbert space  $H_0^1(\Omega)$  endowed with the  $\|\cdot\|_\varepsilon$  norm.

**Definition 3.** Let  $X$  a topological space and consider a closed subset  $A \subset X$ . We say that  $A$  has category  $k$  relative to  $X$  ( $\text{cat}_M A = k$ ) if  $A$  is covered by  $k$  closed sets  $A_j$ ,  $j = 1, \dots, k$ , which are contractible in  $X$ , and  $k$  is the minimum integer with this property. We simply denote  $\text{cat } X = \text{cat}_X X$ .

*Remark 4.* Let  $X_1$  and  $X_2$  be topological spaces. If  $g_1 : X_1 \rightarrow X_2$  and  $g_2 : X_2 \rightarrow X_1$  are continuous operators such that  $g_2 \circ g_1$  is homotopic to the identity on  $X_1$ , then  $\text{cat } X_1 \leq \text{cat } X_2$ .

**Definition 5.** Let  $X$  be any topological space and let  $H_k(X)$  denotes its  $k$ -th homology group with coefficients in  $\mathbb{Q}$ . The Poincaré polynomial  $P_t(X)$  of  $X$  is defined as the following power series in  $t$

$$P_t(X) := \sum_{k \geq 0} (\dim H_k(X)) t^k$$

Actually, if  $X$  is a compact space, we have that  $\dim H_k(X) < \infty$  and this series is finite; in this case,  $P_t(X)$  is a polynomial and not a formal series.

*Remark 6.* Let  $X$  and  $Y$  be topological spaces. If  $f : X \rightarrow Y$  and  $g : Y \rightarrow X$  are continuous operators such that  $g \circ f$  is homotopic to the identity on  $X$ , then  $P_t(Y) = P_t(X) + Z(t)$  where  $Z(t)$  is a polynomial with non-negative coefficients.

These topological tools are classical and can be found, e.g., in [16] and in [5].

### 3. PRELIMINARY RESULTS

Using an idea in a paper of Benci and Fortunato [9] we define the map  $\psi : H_0^1(\Omega) \rightarrow H_0^1(\Omega)$  defined by the equation

$$(2) \quad -\Delta\psi(u) = qu^2 \text{ in } \Omega$$

**Lemma 7.** *The map  $\psi : H_0^1(\Omega) \rightarrow H_0^1(\Omega)$  is of class  $C^2$  with derivatives*

$$(3) \quad \psi'(u)[\varphi] = i^*(2qu\varphi)$$

$$(4) \quad \psi''(u)[\varphi_1, \varphi_2] = i^*(2q\varphi_1\varphi_2)$$

where the operator  $i_\varepsilon^* : L^{p'}, |\cdot|_{\varepsilon, p'} \rightarrow H_\varepsilon$  is the adjoint operator of the immersion operator  $i_\varepsilon : H_\varepsilon \rightarrow L^p, |\cdot|_{\varepsilon, p}$ .

*Proof.* The proof is standard.  $\square$

**Lemma 8.** *The map  $T : H_0^1(\Omega) \rightarrow \mathbb{R}$  given by*

$$T(u) = \int_{\Omega} u^2 \psi(u) dx$$

is a  $C^2$  map and its first derivative is

$$T'(u)[\varphi] = 4 \int_{\Omega} \varphi u \psi(u) dx.$$

*Proof.* The regularity is standard. The first derivative is

$$T'(u)[\varphi] = 2 \int u \varphi \psi(u) + \int u^2 \psi'(u)[\varphi].$$

By (3) and (2) we have

$$\begin{aligned} 2q \int u \varphi \psi(u) &= - \int \Delta(\psi'(u)[\varphi]) \psi(u) = - \int \psi'(u)[\varphi] \Delta\psi(u) = \\ &= \int \psi'(u)[\varphi] qu^2 \end{aligned}$$

and the claim follows.  $\square$

At this point we consider the following functional  $I_\varepsilon \in C^2(H_0^1(\Omega), \mathbb{R})$ .

$$(5) \quad I_\varepsilon(u) = \frac{1}{2} \|u\|_{\varepsilon}^2 + \frac{\omega}{4} G_\varepsilon(u) - \frac{1}{p} |u^+|_{\varepsilon, p}^p$$

where

$$G_\varepsilon(u) = \frac{1}{\varepsilon^3} \int_{\Omega} u^2 \psi(u) dx = \frac{1}{\varepsilon^3} T(u).$$

By Lemma 8 we have

$$I'_\varepsilon(u)[\varphi] = \frac{1}{\varepsilon^3} \int_{\Omega} \varepsilon^2 \nabla u \nabla \varphi + u \varphi + \omega u \psi(u) \varphi - (u^+)^{p-1} \varphi$$

$$I'_\varepsilon(u)[u] = \|u\|_{\varepsilon}^2 + \omega G_\varepsilon(u) - |u^+|_{\varepsilon, p}^p$$

then if  $u$  is a critical points of the functional  $I_\varepsilon$  the pair of positive functions  $(u, \psi(u))$  is a solution of (1).

#### 4. NEHARI MANIFOLD

We define the following Nehari set

$$\mathcal{N}_\varepsilon = \{u \in H_0^1(\Omega) \setminus \{0\} : N_\varepsilon(u) := I'_\varepsilon(u)[u] = 0\}$$

In this section we give an explicit proof of the main properties of the Nehari manifold, although standard, for the sake of completeness

**Lemma 9.**  *$\mathcal{N}_\varepsilon$  is a  $C^2$  manifold and  $\inf_{\mathcal{N}_\varepsilon} \|u\|_\varepsilon > 0$ .*

*Proof.* If  $u \in \mathcal{N}_\varepsilon$ , using that  $N_\varepsilon(u) = 0$ , and  $p > 4$  we have

$$N'_\varepsilon(u)[u] = 2\|u\|_\varepsilon^2 + 4\omega G_\varepsilon(u) - p|u^+|_{\varepsilon,p} = (2-p)\|u\|_\varepsilon + (4-p)\omega G_\varepsilon(u) < 0$$

so  $\mathcal{N}_\varepsilon$  is a  $C^2$  manifold.

We prove the second claim by contradiction. Take a sequence  $\{u_n\}_n \in \mathcal{N}_\varepsilon$  with  $\|u_n\|_\varepsilon \rightarrow 0$  while  $n \rightarrow +\infty$ . Thus, using that  $N_\varepsilon(u) = 0$ ,

$$\|u_n\|_\varepsilon^2 + \omega G_\varepsilon(u_n) = |u_n^+|_{p,\varepsilon}^p \leq C\|u_n\|_\varepsilon^p,$$

so

$$1 < 1 + \frac{\omega G_\varepsilon(u)}{\|u_n\|_\varepsilon} \leq C\|u_n\|_\varepsilon^{p-2} \rightarrow 0$$

and this is a contradiction.  $\square$

*Remark 10.* If  $u \in \mathcal{N}_\varepsilon$ , then

$$\begin{aligned} I_\varepsilon(u) &= \left(\frac{1}{2} - \frac{1}{p}\right)\|u\|_\varepsilon^2 + \omega\left(\frac{1}{4} - \frac{1}{p}\right)G_\varepsilon(u) \\ &= \left(\frac{1}{2} - \frac{1}{p}\right)|u^+|_{p,\varepsilon}^p - \frac{\omega}{4}G_\varepsilon(u) \end{aligned}$$

**Lemma 11.** *It holds Palais-Smale condition for the functional  $I_\varepsilon$  on  $\mathcal{N}_\varepsilon$ .*

*Proof.* We start proving PS condition for  $I_\varepsilon$ . Let  $\{u_n\}_n \in H_0^1(\Omega)$  such that

$$I_\varepsilon(u_n) \rightarrow c \quad |I'_\varepsilon(u_n)[\varphi]| \leq \sigma_n\|\varphi\|_\varepsilon \text{ where } \sigma_n \rightarrow 0$$

We prove that  $\|u_n\|_\varepsilon$  is bounded. Suppose  $\|u_n\|_\varepsilon \rightarrow \infty$ . Then, by PS hypothesis

$$\frac{pI_\varepsilon(u_n) - I'_\varepsilon(u_n)[u_n]}{\|u_n\|_\varepsilon} = \left(\frac{p}{2} - 1\right)\|u_n\|_\varepsilon + \left(\frac{p}{4} - 1\right)\frac{G_\varepsilon(u_n)}{\|u_n\|_\varepsilon} \rightarrow 0$$

and this is a contradiction because  $p > 4$ .

At this point, up to subsequence  $u_n \rightarrow u$  weakly in  $H_0^1(\Omega)$  and strongly in  $L^t(\Omega)$  for each  $2 \leq t < 6$ . Since  $u_n$  is a PS sequence

$$u_n + \omega i_\varepsilon^*(\psi(u_n)u_n) - i_\varepsilon^*((u_n^+)^{p-1}) \rightarrow 0 \text{ in } H_0^1(\Omega)$$

we have only to prove that  $i_\varepsilon^*(\psi(u_n)u_n) \rightarrow i_\varepsilon^*(\psi(u)u)$  in  $H_0^1(\Omega)$ , then we have to prove that

$$\psi(u_n)u_n \rightarrow \psi(u)u \text{ in } L^{t'}$$

We have  $|\psi(u_n)u_n - \psi(u)u|_{\varepsilon,t'} \leq |\psi(u)(u_n - u)|_{\varepsilon,t'} + |(\psi(u_n) - \psi(u))u_n|_{\varepsilon,t'}$ . We get

$$\int_\Omega |\psi(u_n) - \psi(u)|^{\frac{t}{t-1}} |u_n|^{\frac{t}{t-1}} \leq \left(\int_\Omega |\psi(u_n) - \psi(u)|^t\right)^{\frac{1}{t-1}} \left(\int_\Omega |u_n|^{\frac{t}{t-2}}\right)^{\frac{t-2}{t-1}} \rightarrow 0,$$

thus we can conclude easily.

Now we prove PS condition for the constrained functional. Let  $\{u_n\}_n \in \mathcal{N}_\varepsilon$  such that

$$I_\varepsilon(u_n) \rightarrow c \\ |I'_\varepsilon(u_n)[\varphi] - \lambda_n N'(u_n)[\varphi]| \leq \sigma_n \|\varphi\|_\varepsilon \quad \text{with } \sigma_n \rightarrow 0$$

In particular  $I'_\varepsilon(u_n) \left[ \frac{u_n}{\|u_n\|_\varepsilon} \right] - \lambda_n N'(u_n) \left[ \frac{u_n}{\|u_n\|_\varepsilon} \right] \rightarrow 0$ . Then

$$\lambda_n \left\{ (p-2) \|u_n\|_\varepsilon + (p-4) \omega \frac{G_\varepsilon(u_n)}{\|u_n\|_\varepsilon} \right\} \rightarrow 0$$

thus  $\lambda_n \rightarrow 0$  because  $p > 4$ . Since  $N'(u_n) = u_n - i_\varepsilon^*(4\omega\psi(u_n)u_n) - pi_\varepsilon^*(|u_n^+|^{p-1})$  is bounded we obtain that  $\{u_n\}_n$  is a PS sequence for the free functional  $I_\varepsilon$ , and we get the claim  $\square$

**Lemma 12.** *For all  $w \in H_0^1(\Omega)$  such that  $|w^+|_{\varepsilon,p} = 1$  there exists a unique positive number  $t_\varepsilon = t_\varepsilon(w)$  such that  $t_\varepsilon(w)w \in \mathcal{N}_\varepsilon$ .*

*Proof.* We define, for  $t > 0$

$$H(t) = I_\varepsilon(tw) = \frac{1}{2}t^2\|w\|_\varepsilon^2 + \frac{t^4}{4}\omega G_\varepsilon(w) - \frac{t^p}{p}.$$

Thus

$$(6) \quad H'(t) = t (\|w\|_\varepsilon^2 + t^2\omega G_\varepsilon(w) - t^{p-2})$$

$$(7) \quad H''(t) = \|w\|_\varepsilon^2 + 3t^2\omega G_\varepsilon(w) - (p-1)t^{p-2}$$

By (6) there exists  $t_\varepsilon > 0$  such that  $H'(t_\varepsilon)$ . Moreover, by (6), (7) and because  $p > 4$  we have that  $H''(t_\varepsilon) < 0$ , so  $t_\varepsilon$  is unique.  $\square$

## 5. MAIN INGREDIENT OF THE PROOF

We sketch the proof of Theorem 1. First of all, since the functional  $I_\varepsilon \in C^2$  is bounded below and satisfies PS condition on the complete  $C^2$  manifold  $\mathcal{N}_\varepsilon$ , we have, by well known results, that  $I_\varepsilon$  has at least  $\text{cat } I_\varepsilon^d$  critical points in the sublevel

$$I_\varepsilon^d = \{u \in H^1 : I_\varepsilon(u) \leq d\}.$$

We prove that, for  $\varepsilon$  and  $\delta$  small enough, it holds

$$\text{cat } \Omega \leq \text{cat } (\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta})$$

where

$$m_\infty := \inf_{\mathcal{N}_\infty} \frac{1}{2} \int_{\mathbb{R}^3} |\nabla v|^2 + v^2 dx - \frac{1}{p} \int_{\mathbb{R}^3} |v|^p dx$$

$$\mathcal{N}_\infty = \left\{ v \in H^1(\mathbb{R}^3) \setminus \{0\} : \int_{\mathbb{R}^3} |\nabla v|^2 + v^2 dx = \int_{\mathbb{R}^3} |v|^p dx \right\}.$$

To get the inequality  $\text{cat } \Omega \leq \text{cat } (\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta})$  we build two continuous operators

$$\begin{aligned} \Phi_\varepsilon & : \Omega^- \rightarrow \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta} \\ \beta & : \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta} \rightarrow \Omega^+ \end{aligned}$$

where

$$\Omega^- = \{x \in \Omega : d(x, \partial\Omega) < r\}$$

$$\Omega^+ = \{x \in \mathbb{R}^3 : d(x, \partial\Omega) < r\}$$

with  $r$  small enough so that  $\text{cat}(\Omega^-) = \text{cat}(\Omega^+) = \text{cat}(\Omega)$ .

Following an idea in [7], we build these operators  $\Phi_\varepsilon$  and  $\beta$  such that  $\beta \circ \Phi_\varepsilon : \Omega^- \rightarrow \Omega^+$  is homotopic to the immersion  $i : \Omega^- \rightarrow \Omega^+$ . By the properties of Lusternik Schinerlmann category we have

$$\text{cat } \Omega \leq \text{cat} (\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta})$$

which ends the proof of Theorem 1.

Concerning Theorem 2, we can re-state classical results contained in [5, 8] in the following form.

**Theorem 13.** *Let  $I_\varepsilon$  be the functional (5) on  $H^1(\Omega)$  and let  $K_\varepsilon$  be the set of its critical points. If all its critical points are non-degenerate then*

$$(8) \quad \sum_{u \in K_\varepsilon} t^{\mu(u)} = tP_t(\Omega) + t^2(P_t(\Omega) - 1) + t(1+t)Q(t)$$

where  $Q(t)$  is a polynomial with non-negative integer coefficients and  $\mu(u)$  is the Morse index of the critical point  $u$ .

By Remark 6 and by means of the maps  $\Phi_\varepsilon$  and  $\beta$  we have that

$$(9) \quad P_t(\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta}) = P_t(\Omega) + Z(t)$$

where  $Z(t)$  is a polynomial with non-negative coefficients. Provided that  $\inf_\varepsilon m_\varepsilon =: \alpha > 0$ , because  $\lim_{\varepsilon \rightarrow 0} m_\varepsilon = m_\infty$  (see 20), we have the following relations [5, 8]

$$(10) \quad P_t(I_\varepsilon^{m_\infty + \delta}, I_\varepsilon^{\alpha/2}) = tP_t(\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta})$$

$$(11) \quad P_t(H_0^1(\Omega), I_\varepsilon^{m_\infty + \delta}) = t(P_t(I_\varepsilon^{m_\infty + \delta}, I_\varepsilon^{\alpha/2}) - t)$$

$$(12) \quad \sum_{u \in K_\varepsilon} t^{\mu(u)} = P_t(H_0^1(\Omega), I_\varepsilon^{m_\infty + \delta}) + P_t(I_\varepsilon^{m_\infty + \delta}, I_\varepsilon^{\alpha/2}) + (1+t)\tilde{Q}(t)$$

where  $\tilde{Q}(t)$  is a polynomial with non-negative integer coefficients. Hence, by (9), (10), (11), (12) we obtain (8). At this point, evaluating equation (8) for  $t = 1$  we obtain the claim of Theorem 2

## 6. THE MAP $\Phi_\varepsilon$

For every  $\xi \in \Omega^-$  we define the function

$$W_{\xi, \varepsilon}(x) = U_\varepsilon(x - \xi)\chi(|x - \xi|)$$

where  $\chi : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  where  $\chi \equiv 1$  for  $t \in [0, r/2)$ ,  $\chi \equiv 0$  for  $t > r$  and  $|\chi'(t)| \leq 2/r$ .

We can define a map

$$\begin{aligned} \Phi_\varepsilon &: \Omega^- \rightarrow \mathcal{N}_\varepsilon \\ \Phi_\varepsilon(\xi) &= t_\varepsilon(W_{\xi, \varepsilon})W_{\xi, \varepsilon} \end{aligned}$$

*Remark 14.* We have that the following limits hold uniformly with respect to  $\xi \in \Omega$

$$\begin{aligned} \|W_{\varepsilon, \xi}\|_\varepsilon &\rightarrow \|U\|_{H^1(\mathbb{R}^3)} \\ |W_{\varepsilon, \xi}|_{\varepsilon, t} &\rightarrow \|U\|_{L^t(\mathbb{R}^3)} \text{ for all } 2 \leq t \leq 6 \end{aligned}$$

**Lemma 15.** *There exists  $\bar{\varepsilon} > 0$  and a constant  $c > 0$  such that*

$$G_\varepsilon(W_{\varepsilon, \xi}) = \frac{1}{\varepsilon^3} \int_{\Omega} qW_{\varepsilon, \xi}^2(x)\psi(W_{\varepsilon, \xi})dx < c\varepsilon^2$$

*Proof.* It holds

$$\begin{aligned} \|\psi(W_{\varepsilon,\xi})\|_{H_0^1(\Omega)}^2 &= \int_{\Omega} q W_{\varepsilon,\xi}^2(x) \psi(W_{\varepsilon,\xi}) dx \leq q \|\psi(W_{\varepsilon,\xi})\|_{L^6(\Omega)} \left( \int_{\Omega} W_{\varepsilon,\xi}^{12/5} dx \right)^{5/6} \\ &\leq c \|\psi(W_{\varepsilon,\xi})\|_{H_0^1(\Omega)} \left( \frac{1}{\varepsilon^3} \int_{\Omega} W_{\varepsilon,\xi}^{12/5} dx \right)^{5/6} \varepsilon^{5/2} \end{aligned}$$

By Remark 14 we have that  $\|\psi(W_{\varepsilon,\xi})\|_{H_0^1(\Omega)} \leq \varepsilon^{5/2}$  and the claim follows by applying again Cauchy Schwartz inequality.  $\square$

**Proposition 16.** *For all  $\varepsilon > 0$  the map  $\Phi_{\varepsilon}$  is continuous. Moreover for any  $\delta > 0$  there exists  $\varepsilon_0 = \varepsilon_0(\delta)$  such that, if  $\varepsilon < \varepsilon_0$  then  $I_{\varepsilon}(\Phi_{\varepsilon}(\xi)) < m_{\infty} + \delta$ .*

*Proof.* It is easy to see that  $\Phi_{\varepsilon}$  is continuous because  $t_{\varepsilon}(w)$  depends continuously on  $w \in H_0^1$ .

At this point we prove that  $t_{\varepsilon}(W_{\varepsilon,\xi}) \rightarrow 1$  uniformly with respect to  $\xi \in \Omega$ . In fact, by Lemma 12  $t_{\varepsilon}(W_{\varepsilon,\xi})$  is the unique solution of

$$\|W_{\varepsilon,\xi}\|_{\varepsilon}^2 + t^2 \omega G_{\varepsilon}(W_{\varepsilon,\xi}) - t^{p-2} |W_{\varepsilon,\xi}|_{\varepsilon,p}^p = 0.$$

By Remark 14 and Lemma 15 we have the claim.

Now, we have

$$I_{\varepsilon}(t_{\varepsilon}(W_{\varepsilon,\xi})W_{\varepsilon,\xi}) = \left( \frac{1}{2} - \frac{1}{p} \right) \|W_{\varepsilon,\xi}\|_{\varepsilon}^2 t_{\varepsilon}^2 + \omega \left( \frac{1}{4} - \frac{1}{p} \right) t_{\varepsilon}^4 G_{\varepsilon}(W_{\varepsilon,\xi})$$

Again, by Remark 14 and Lemma 15 we have

$$I_{\varepsilon}(t_{\varepsilon}(W_{\varepsilon,\xi})W_{\varepsilon,\xi}) \rightarrow \left( \frac{1}{2} - \frac{1}{p} \right) \|U\|_{H^1(\mathbb{R}^3)}^2 = m_{\infty}$$

that concludes the proof.  $\square$

*Remark 17.* We set

$$m_{\varepsilon} = \inf_{\mathcal{N}_{\varepsilon}} I_{\varepsilon}.$$

By Proposition 16 we have that

$$(13) \quad \limsup_{\varepsilon \rightarrow 0} m_{\varepsilon} \leq m_{\infty}.$$

## 7. THE MAP $\beta$

For any  $u \in \mathcal{N}_{\varepsilon}$  we can define a point  $\beta(u) \in \mathbb{R}^3$  by

$$\beta(u) = \frac{\int_{\Omega} x |u^+|^p dx}{\int_{\Omega} |u^+|^p dx}.$$

The function  $\beta$  is well defined in  $\mathcal{N}_{\varepsilon}$  because, if  $u \in \mathcal{N}_{\varepsilon}$ , then  $u^+ \neq 0$ .

We have to prove that, if  $u \in \mathcal{N}_{\varepsilon} \cap I_{\varepsilon}^{m_{\infty}+\delta}$  then  $\beta(u) \in \Omega^+$ .

Let us consider partitions of  $\Omega$ . For a given  $\varepsilon > 0$  we say that a finite partition  $\mathcal{P}_{\varepsilon} = \{P_j^{\varepsilon}\}_{j \in \Lambda_{\varepsilon}}$  of  $\Omega$  is a “good” partition if: for any  $j \in \Lambda_{\varepsilon}$  the set  $P_j^{\varepsilon}$  is closed;  $P_i^{\varepsilon} \cap P_j^{\varepsilon} \subset \partial P_i^{\varepsilon} \cap \partial P_j^{\varepsilon}$  for any  $i \neq j$ ; there exist  $r_1(\varepsilon), r_2(\varepsilon) > 0$  such that there are points  $q_j^{\varepsilon} \in P_j^{\varepsilon}$  for which  $B(q_j^{\varepsilon}, \varepsilon) \subset P_j^{\varepsilon} \subset B(q_j^{\varepsilon}, r_2(\varepsilon)) \subset B_g(q_j^{\varepsilon}, r_1(\varepsilon))$ , with  $r_1(\varepsilon) \geq r_2(\varepsilon) \geq C\varepsilon$  for some positive constant  $C$ ; lastly, there exists a finite number

$\nu \in \mathbb{N}$  such that every  $x \in \Omega$  is contained in at most  $\nu$  balls  $B(q_j^\varepsilon, r_1(\varepsilon))$ , where  $\nu$  does not depends on  $\varepsilon$ .

**Lemma 18.** *There exists a constant  $\gamma > 0$  such that, for any  $\delta > 0$  and for any  $\varepsilon < \varepsilon_0(\delta)$  as in Proposition 16, given any “good” partition  $\mathcal{P}_\varepsilon = \{P_j^\varepsilon\}_j$  of the domain  $\Omega$  and for any function  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta}$  there exists, for an index  $\bar{j}$  a set  $P_{\bar{j}}^\varepsilon$  such that*

$$\frac{1}{\varepsilon^3} \int_{P_{\bar{j}}^\varepsilon} |u^+|^p dx \geq \gamma.$$

*Proof.* Taking in account that  $I'(u)[u] = 0$  we have

$$\begin{aligned} \|u\|_\varepsilon^2 &= |u^+|_{\varepsilon,p}^p - \frac{1}{\varepsilon^3} \int_{\Omega} \omega u^2 \psi(u) \leq |u^+|_{\varepsilon,p}^p = \sum_j \frac{1}{\varepsilon^3} \int_{P_j} |u^+|^p \\ &= \sum_j |u_j^+|_{\varepsilon,p}^p = \sum_j |u_j^+|_{\varepsilon,p}^{p-2} |u_j^+|_{\varepsilon,p}^2 \leq \max_j \{|u_j^+|_{\varepsilon,p}^{p-2}\} \sum_j |u_j^+|_{\varepsilon,p}^2 \end{aligned}$$

where  $u_j^+$  is the restriction of the function  $u^+$  on the set  $P_j$ .

At this point, arguing as in [6, Lemma 5.3], we prove that there exists a constant  $C > 0$  such that

$$\sum_j |u_j^+|_{\varepsilon,p}^2 \leq C\nu \|u^+\|_\varepsilon^2,$$

thus

$$\max_j \{|u_j^+|_{\varepsilon,p}^{p-2}\} \geq \frac{1}{C\nu}$$

that concludes the proof.  $\square$

**Proposition 19.** *For any  $\eta \in (0, 1)$  there exists  $\delta_0 < m_\infty$  such that for any  $\delta \in (0, \delta_0)$  and any  $\varepsilon \in (0, \varepsilon_0(\delta))$  as in Proposition 16, for any function  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + \delta}$  we can find a point  $q = q(u) \in \Omega$  such that*

$$\frac{1}{\varepsilon^3} \int_{B(q, r/2)} (u^+)^p > (1 - \eta) \frac{2p}{p-2} m_\infty.$$

*Proof.* First, we prove the proposition for  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty + 2\delta}$ .

By contradiction, we assume that there exists  $\eta \in (0, 1)$  such that we can find two sequences of vanishing real number  $\delta_k$  and  $\varepsilon_k$  and a sequence of functions  $\{u_k\}_k$  such that  $u_k \in \mathcal{N}_{\varepsilon_k}$ ,

(14)

$$m_{\varepsilon_k} \leq I_{\varepsilon_k}(u_k) = \left( \frac{1}{2} - \frac{1}{p} \right) \|u_k\|_{\varepsilon_k}^2 + \omega \left( \frac{1}{4} - \frac{1}{p} \right) G_{\varepsilon_k}(u_k) \leq m_{\varepsilon_k} + 2\delta_k \leq m_\infty + 3\delta_k$$

for  $k$  large enough (see Remark 17), and, for any  $q \in \Omega$ ,

$$\frac{1}{\varepsilon_k^3} \int_{B(q, r/2)} (u_k^+)^p \leq (1 - \eta) \frac{2p}{p-2} m_\infty.$$

By Ekeland principle and by definition of  $\mathcal{N}_{\varepsilon_k}$  we can assume

$$(15) \quad |I'_{\varepsilon_k}(u_k)[\varphi]| \leq \sigma_k \|\varphi\|_{\varepsilon_k} \text{ where } \sigma_k \rightarrow 0.$$

By Lemma 18 there exists a set  $P_k^{\varepsilon_k} \in \mathcal{P}_{\varepsilon_k}$  such that

$$\frac{1}{\varepsilon_k^3} \int_{P_k^{\varepsilon_k}} |u_k^+|^p dx \geq \gamma.$$

We choose a point  $q_k \in P_k^{\varepsilon_k}$  and we define, for  $z \in \Omega_{\varepsilon_k} := \frac{1}{\varepsilon_k}(\Omega - q_k)$

$$w_k(z) = u_k(\varepsilon_k z + q_k) = u_k(x).$$

We have that  $w_k \in H_0^1(\Omega_{\varepsilon_k}) \subset H^1(\mathbb{R}^3)$ . By equation (14) we have

$$\|w_k\|_{H^1(\mathbb{R}^3)}^2 = \|u_k\|_{\varepsilon_k}^2 \leq C.$$

So  $w_k \rightarrow w$  weakly in  $H^1(\mathbb{R}^3)$  and strongly in  $L_{\text{loc}}^t(\mathbb{R}^3)$ .

We set  $\psi(u_k)(x) := \psi_k(x) = \psi_k(\varepsilon_k z + q_k) := \tilde{\psi}_k(z)$  where  $x \in \Omega$  and  $z \in \Omega_{\varepsilon_k}$ . It is easy to verify that

$$-\Delta_z \tilde{\psi}_k(z) = \varepsilon_k^2 q w_k^2(z).$$

With abuse of language we set

$$\tilde{\psi}_k(z) = \psi(\varepsilon_k w_k).$$

Thus

$$\begin{aligned} I_{\varepsilon_k}(u_k) &= \frac{1}{2}\|u_k\|_{\varepsilon_k}^2 - \frac{1}{p}|u_k^+|_{\varepsilon_k,p}^p + \frac{\omega}{4}\frac{1}{\varepsilon_k^3} \int_{\Omega} q u_k^2 \psi(u_k) = \\ (16) \quad &= \frac{1}{2}\|w_k\|_{H^1(\mathbb{R}^3)}^2 - \frac{1}{p}\|w_k^+\|_{L^p(\mathbb{R}^3)}^p + \frac{\omega}{4} \int_{\Omega_{\varepsilon_k}} q w_k^2 \psi(\varepsilon_k w_k) = \\ &= \frac{1}{2}\|w_k\|_{H^1(\mathbb{R}^3)}^2 - \frac{1}{p}\|w_k^+\|_{L^p(\mathbb{R}^3)}^p + \varepsilon_k^2 \frac{\omega}{4} \int_{\mathbb{R}^3} q w_k^2 \psi(w_k) := E_{\varepsilon_k}(w_k) \end{aligned}$$

By definition of  $E_{\varepsilon_k} : H^1(\mathbb{R}^3) \rightarrow \mathbb{R}$ , we get  $E_{\varepsilon_k}(w_k) \rightarrow m_{\infty}$ .

Given any  $\varphi \in C_0^\infty(\mathbb{R}^3)$  we set  $\varphi(x) = \varphi(\varepsilon_k z + q_k) := \tilde{\varphi}_k(z)$ . For  $k$  large enough we have that  $\text{supp } \tilde{\varphi}_k \subset \Omega$  and, by (15), that  $E'_{\varepsilon_k}(w_k)[\varphi] = I'_{\varepsilon_k}(u_k)[\tilde{\varphi}_k] \rightarrow 0$ . Moreover, by definiton of  $E_{\varepsilon_k}$  and by Lemma 8 we have

$$\begin{aligned} E'_{\varepsilon_k}(w_k)[\varphi] &= \langle w_k, \varphi \rangle_{H^1(\mathbb{R}^3)} - \int_{\mathbb{R}^3} |w_k^+|^{p-1} \varphi + \omega \varepsilon_k^2 \int_{\mathbb{R}^3} q w_k \psi(w_k) \varphi + \\ &\rightarrow \langle w, \varphi \rangle_{H^1(\mathbb{R}^3)} - \int_{\mathbb{R}^3} |w^+|^{p-1} \varphi. \end{aligned}$$

Thus  $w$  is a weak solution of

$$-\Delta w + w = (w^+)^{p-1} \text{ on } \mathbb{R}^3.$$

By Lemma 18 and by the choice of  $q_k$  we have that  $w \neq 0$ , so  $w > 0$ .

Arguing as in (16), and using that  $u_k \in \mathcal{N}_{\varepsilon_k}$  we have

$$\begin{aligned} (17) \quad I_{\varepsilon_k}(u_k) &= \left(\frac{1}{2} - \frac{1}{p}\right)\|u_k\|_{\varepsilon_k}^2 + \omega \left(\frac{1}{4} - \frac{1}{p}\right) \frac{1}{\varepsilon_k^3} \int_{\Omega} q u_k^2 \psi(u_k) \\ &= \left(\frac{1}{2} - \frac{1}{p}\right)\|w_k\|_{H^1(\mathbb{R}^3)}^2 + \varepsilon_k^2 \omega \left(\frac{1}{4} - \frac{1}{p}\right) \int_{\mathbb{R}^3} q w_k^2 \psi(w_k) \rightarrow m_{\infty} \end{aligned}$$

and

$$\begin{aligned} (18) \quad I_{\varepsilon_k}(u_k) &= \left(\frac{1}{2} - \frac{1}{p}\right)|u_k^+|_{p,\varepsilon_k}^p - \frac{\omega}{4}\frac{1}{\varepsilon_k^3} \int_{\Omega} q u_k^2 \psi(u_k) \\ &= \left(\frac{1}{2} - \frac{1}{p}\right)|w_k^+|_p^p - \varepsilon_k^2 \frac{\omega}{4} \int_{\mathbb{R}^3} q w_k^2 \psi(w_k) \rightarrow m_{\infty}. \end{aligned}$$

So, by (17) we have that  $\|w\|_{H^1(\mathbb{R}^3)}^2 = \frac{2p}{p-2}m_{\infty}$  and that  $\left(\frac{1}{2} - \frac{1}{p}\right)\|w_k\|_{H^1(\mathbb{R}^3)}^2 \rightarrow m_{\infty}$  and we conclude that  $w_k \rightarrow w$  strongly in  $H^1(\mathbb{R}^3)$ .

Given  $T > 0$ , by the definiton of  $w_k$  we get, for  $k$  large enough

$$\begin{aligned} |w_k^+|_{L^p(B(0,T))}^p &= \frac{1}{\varepsilon_k^3} \int_{B(q_k, \varepsilon_k T)} |u_k^+|^p dx \leq \frac{1}{\varepsilon_k^3} \int_{B(q_k, r/2)} |u_k^+|^p dx \\ (19) \quad &\leq (1-\eta) \frac{2p}{p-2} m_\infty. \end{aligned}$$

Then we have the contradiction. In fact, by (18) we have  $\left(\frac{1}{2} - \frac{1}{p}\right) |w_k^+|_p^p \rightarrow m_\infty$  and this contradicts (19). At this point we have proved the claim for  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+2\delta}$ . Now, by the thesis for  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+2\delta}$  and by (18) we have

$$I_{\varepsilon_k}(u_k) = \left(\frac{1}{2} - \frac{1}{p}\right) |u_k^+|_{p,\varepsilon_k}^p + O(\varepsilon^2) \geq (1-\eta)m_\infty + O(\varepsilon^2)$$

and, passing to the limit,

$$\liminf_{k \rightarrow \infty} m_{\varepsilon_k} \geq m_\infty.$$

This, combined by (13) gives us that

$$(20) \quad \lim_{\varepsilon \rightarrow 0} m_\varepsilon = m_\infty.$$

Hence, when  $\varepsilon, \delta$  are small enough,  $\mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta} \subset \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+2\delta}$  and the general claim follows.  $\square$

**Proposition 20.** *There exists  $\delta_0 \in (0, m_\infty)$  such that for any  $\delta \in (0, \delta_0)$  and any  $\varepsilon \in (0, \varepsilon(\delta_0))$  (see Proposition 16), for every function  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta}$  it holds  $\beta(u) \in \Omega^+$ . Moreover the composition*

$$\beta \circ \Phi_\varepsilon : \Omega^- \rightarrow \Omega^+$$

is s homotopic to the immersion  $i : \Omega^- \rightarrow \Omega^+$

*Proof.* By Proposition 19, for any function  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta}$ , for any  $\eta \in (0, 1)$  and for  $\varepsilon, \delta$  small enough, we can find a point  $q = q(u) \in \Omega$  such that

$$\frac{1}{\varepsilon^3} \int_{B(q, r/2)} (u^+)^p > (1-\eta) \frac{2p}{p-2} m_\infty.$$

Moreover, since  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta}$  we have

$$I_\varepsilon(u) = \left(\frac{p-2}{2p}\right) |u^+|_{p,\varepsilon}^p - \frac{\omega}{4} \frac{1}{\varepsilon^3} \int_{\Omega} qu^2 \psi(u) \leq m_\infty + \delta.$$

Now, arguing as in Lemma 15 we have that

$$\|\psi(u)\|_{H^1(\Omega)}^2 = q \int_{\Omega} \psi(u) u^2 \leq C \|\psi(u)\|_{H^1(\Omega)} \left( \int_{\Omega} u^{12/5} \right)^{5/6},$$

so  $\|\psi(u)\|_{H^1(\Omega)} \leq \left( \int_{\Omega} u^{12/5} \right)^{5/6}$ , then

$$\begin{aligned} \frac{1}{\varepsilon^3} \int \psi(u) u^2 &\leq \frac{1}{\varepsilon^3} \|\psi\|_{H^1(\Omega)} \left( \int_{\Omega} u^{12/5} \right)^{5/6} \leq C \frac{1}{\varepsilon^3} \left( \int_{\Omega} u^{12/5} \right)^{5/3} \\ &\leq C \varepsilon^2 |u|_{12/5, \varepsilon}^4 \leq C \varepsilon^2 \|u\|_{\varepsilon}^4 \leq C \varepsilon^2 \end{aligned}$$

because  $\|u\|_{\varepsilon}$  is bounded since  $u \in \mathcal{N}_\varepsilon \cap I_\varepsilon^{m_\infty+\delta}$ .

Hence, provided we choose  $\varepsilon(\delta_0)$  small enough, we have

$$\left( \frac{p-2}{2p} \right) |u^+|_{p,\varepsilon}^p \leq m_\infty + 2\delta_0.$$

So,

$$\frac{\frac{1}{\varepsilon^3} \int_{B(q,r/2)} (u^+)^p}{|u^+|_{p,\varepsilon}^p} > \frac{1-\eta}{1+2\delta_0/m_\infty}$$

Finally,

$$\begin{aligned} |\beta(u) - q| &\leq \frac{\left| \frac{1}{\varepsilon^3} \int_{\Omega} (x-q)(u^+)^p \right|}{|u^+|_{p,\varepsilon}^p} \\ &\leq \frac{\left| \frac{1}{\varepsilon^3} \int_{B(q,r/2)} (x-q)(u^+)^p \right|}{|u^+|_{p,\varepsilon}^p} + \frac{\left| \frac{1}{\varepsilon^3} \int_{\Omega \setminus B(q,r/2)} (x-q)(u^+)^p \right|}{|u^+|_{p,\varepsilon}^p} \\ &\leq \frac{r}{2} + 2\text{diam}(\Omega) \left( 1 - \frac{1-\eta}{1+2\delta_0/m_\infty} \right), \end{aligned}$$

so, choosing  $\eta$ ,  $\delta_0$  and  $\varepsilon(\delta_0)$  small enough we proved the first claim. The second claim is standard.  $\square$

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